# ON THE ELLIPTIC STARK CONJECTURE AT PRIMES OF MULTIPLICATIVE REDUCTION

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ABSTRACT. In [DLR], Darmon, Lauder and Rotger formulated a p-adic elliptic Stark conjecture for the twist of an elliptic curve  $E/\mathbb{Q}$  by the self-dual tensor product  $\rho_1\otimes\rho_2$  of two odd and two-dimensional Artin representations. The authors of loc. cit. provided abundant numerical evidence and proved the conjecture in the special setting where p is a prime of good reduction for E and  $\rho_1$  and  $\rho_2$  are induced from finite order characters  $\psi_g$ ,  $\psi_h$  of the same imaginary quadratic field. The key step in their proof is a factorization of one-variable p-adic L-functions, where  $\psi_g$  varies in a p-adic family of Hecke characters.

The main goal of this article is to prove a new case of the conjecture, placing ourselves in the setting where p is a prime of multiplicative reduction for E. In order to achieve our theorem, we need to work with two-variable p-adic L-functions, where the weight 2 cusp form associated with E also moves independently along a Hida family. Our main result then follows from a factorization of p-adic L-series extending to two-variables the one obtained in [DLR]. On the way we also generalize to our setting the results obtained in [CR].

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## 1. Introduction

Fix a prime number  $p \geq 3$  and three positive integers  $N_f, N_g, N_h$  such that  $p \nmid N_f N_g N_h$ . Set  $N = \text{lcm}(N_f, N_g, N_h)$  and let  $\chi : (\mathbb{Z}/N\mathbb{Z})^{\times} \to \mathbb{C}^{\times}$  be a Dirichlet character. Let

$$f \in S_2(pN_f), \quad g \in M_1(N_q, \bar{\chi}), \quad h \in M_1(N_h, \chi)$$

be a triple of newforms of weights (2,1,1), levels  $(pN_f,N_g,N_h)$  and nebentype characters  $(1,\bar{\chi},\chi)$ .

Assume for simplicity that the Fourier coefficients of f are rational and let  $E/\mathbb{Q}$  be an elliptic curve in the isogeny class associated with f by Eichler-Shimura. Let also  $\varrho_g$ ,  $\varrho_h$  denote the odd, two-dimensional Artin representations associated by Deligne-Serre to g and h. This paper focuses on the arithmetic of the elliptic curve E twisted by the four-dimensional Artin representation

$$\varrho_{gh} := \varrho_g \otimes \varrho_h,$$

which is self-dual because the nebentype characters of g and h are inverse to each other.

Let  $L(E, \varrho_{gh}, s)$  denote the Hasse-Weil-Artin L-series associated with the compatible system of Galois representations afforded by the twist of E by  $\varrho_{gh}$ . A direct consequence of the

Eichler-Shimura and Deligne-Serre constructions is the identity

$$L(E, \varrho_{gh}, s) = L(f, g, h, s),$$

which recasts the L-function associated with the pair  $(E, \varrho_{gh})$  as Garrett's [Gar] triple-product L-function attached to the triple of eigenforms (f, g, h).

Assuming that  $(pN_f, N_gN_h) = 1$ , the global sign of the functional equation of  $L(E, \varrho_{gh}, s)$  is +1. We refer e.g. to [Pr90] for a proof of this simple fact and for similar assertions under more relaxed hypotheses.

Assume further that  $L(E, \varrho_{gh}, s)$  vanishes at the central critical point s = 1. It follows that the order of vanishing of  $L(E, \varrho_{gh}, s)$  at s = 1 is even and hence at least 2.

In [DLR], Darmon, Lauder and one of the authors of this note introduced certain *p-adic* iterated integrals attached to the triple (f, g, h), which they view as *p*-adic avatars of the leading term of  $L(E, \varrho_{gh}, s)$  at s = 1. These iterated integrals are defined as linear forms

$$S_2(Np)_L[f] \times M_1(Np,\chi)_L^{\vee}[g_{\alpha}] \times M_1(Np,\chi)_L[h] \longrightarrow \mathbb{C}_p$$

$$(\check{f},\check{\gamma},\check{h}) \mapsto \int_{\check{\gamma}} \check{f} \cdot \check{h}$$

on the space of triples of test vectors of level Np on which the good Hecke operators act with the same system of eigenvalues as f,  $g_{\alpha}$  and h respectively. Here  $g_{\alpha}$  is a choice of (necessarily ordinary) p-stabilization of g, on which  $U_p$  acts with eigenvalue  $\alpha$ .

Under our hypotheses and an additional assumption referred to as Hypothesis C in [DLR] on the local geometry of the points corresponding to the ordinary stabilisations of g in the eigencurve, the authors of loc. cit. propose a conjectural formula for these iterated integrals which may be regarded as a p-adic avatar of a putative Gross-Zagier formula in rank 2.

More precisely, these iterated integrals are expected to vanish systematically when the second derivative  $L''(E, \varrho_{gh}, 1)$  vanishes; if instead  $L''(E, \varrho_{gh}, 1) \neq 0$ , the main conjecture of [DLR] predicts that the values of the iterated integrals attached to (f, g, h) encode the logarithm of a suitable Gross-Stark unit and a  $2 \times 2$  p-adic regulator involving the formal group logarithms of certain Stark points in the Mordell-Weil group of E over the number field cut out by  $\varrho_{gh}$ . This conjecture was christened in [DLR] as the *elliptic Stark conjecture*.

The main theorem of [DLR, Ch. 2 & 3] shows that the elliptic Stark conjecture holds when p is a prime of good reduction for E, while g and h are binary theta series attached to a common imaginary quadratic field K in which p splits. This is done by recasting both sides of the conjectural formula, that is to say p-adic iterated integrals on one side and elliptic units and Heegner points on the other, as values of various p-adic L-functions at points lying outside their region of interpolation.

The first main theorem of this article extends the above result to the case where p is a prime of multiplicative reduction of E. In order to state the result precisely, let us briefly review the ingredients appearing in its formulation.

Fix throughout embeddings  $\bar{\mathbb{Q}} \subset \mathbb{C}$  and  $\bar{\mathbb{Q}} \subset \mathbb{C}_p$  so that every number field shall be regarded without further mention as a subfield of both  $\mathbb{C}$  and  $\mathbb{C}_p$ .

Let K be an imaginary quadratic field of discriminant  $-D_K$  and let us denote by  $\chi_K$  the quadratic Dirichlet character associated with the extension  $K/\mathbb{Q}$ . Assume that p splits in K, i.e.  $\chi_K(p) = 1$ .

Let g and h be theta series associated with two finite order characters  $\psi_g, \psi_h : G_K \longrightarrow L^{\times}$  of K of conductors  $\mathfrak{c}_g$  and  $\mathfrak{c}_h \subseteq \mathcal{O}_K$  respectively, with values in a finite extension L of  $\mathbb{Q}$ . They are weight one forms of levels  $D_K \mathrm{N}_{K/\mathbb{Q}}(\mathfrak{c}_g)$  and  $D_K \mathrm{N}_{K/\mathbb{Q}}(\mathfrak{c}_h)$ , having nebentype characters  $\chi_g$  and  $\chi_h$  respectively.

Define  $\psi'_g: G_K \longrightarrow L^{\times}$  to be the  $\operatorname{Gal}(K/\mathbb{Q})$ -conjugate of the character  $\psi_g$ , namely the character given by the rule  $\psi'_g(\sigma) = \psi_g(\tau \sigma \tau^{-1})$  for any  $\sigma \in G_K$  and any  $\tau \in G_{\mathbb{Q}} \setminus G_K$ .

We assume that q satisfies the classicality hypotheses of [DLR] at the prime p. As explained in loc. cit, this is automatically satisfied

- when  $\psi_g/\psi_g'=1$  (i.e. g is Eisenstein); when  $\psi_g/\psi_g'\neq 1$  (i.e. g is cuspidal), if we impose  $\psi(\wp)\neq \psi'(\wp)$  and that  $\rho_g$  is not induced from a character of a real quadratic field in which p splits.

As above, we assume that  $\chi_h = \chi_q^{-1} (= \chi)$ , which implies that the characters

$$\psi_1 = \psi_g \psi_h$$
 and  $\psi_2 = \psi_g \psi_h'$ 

are ring class characters of K, associated with orders  $\mathcal{O}_{c_1}$  and  $\mathcal{O}_{c_2}$  in  $\mathcal{O}_K$  of conductors  $c_1$ and  $c_2$  respectively. Let H denote the ring class field associated with the order  $\mathcal{O}_c$  in K of conductor  $c := \text{lcm}(c_1, c_2)$  of K, which contains the fields cut out by  $\psi_1$  and  $\psi_2$ .

A simple exercise shows that the Artin representation  $\varrho_{gh} = \varrho_g \otimes \varrho_h$  decomposes as

$$\varrho_{gh} = \rho_1 \oplus \rho_2 = V_{\psi_1} \oplus V_{\psi_2},$$

and the Artin formalism in turn implies that there is a factorisation of L-series

(1) 
$$L(E, \varrho_{qh}, s) = L(E, \rho_1, s) \cdot L(E, \rho_2, s) = L(E/K, \psi_1, s) \cdot L(E/K, \psi_2, s).$$

Recall that  $N = \text{lcm}(N_f, N_g, N_h) = \text{lcm}(N_f, D_K N_{K/\mathbb{Q}}(\mathfrak{c}_g), D_K N_{K/\mathbb{Q}}(\mathfrak{c}_h))$ . Assume in this introduction that N is square-free,  $(N_f D_K, c) = 1$  and the following assumption holds (which is exactly the one imposed in  $[BDP13, \S4.1]$ ):

**Heegner hypothesis:** All prime divisors of  $N_f$  split or ramify in K, and the local signs at primes dividing  $gcd(N_f, D_K)$  of  $L(E, \rho_1, s)$  and  $L(E, \rho_2, s)$  are +1.

We refer to [Gro, §11 and 22] for more details about this hypothesis. As shown in loc. cit., this assumption automatically implies that the local signs of  $L(E, \rho_1, s)$  and  $L(E, \rho_2, s)$  are +1 at all finite primes. Since the local sign at  $\infty$  of  $L(E, \rho_i, s)$  is -1, this assumption implies that the global signs of the functional equations satisfied by  $L(E/K, \psi_1, s)$  and  $L(E/K, \psi_2, s)$ are -1, hence the analytic ranks  $r_{\rm an}(E, \varrho_{\psi_1})$  and  $r_{\rm an}(E, \varrho_{\psi_2})$  are both odd.

It also follows from (1) that  $L(E, \varrho_{gh}, 1) = 0$  and thus Hypothesis B of [DLR] is satisfied. Moreover, since the local sign at any place of  $L(E, \varrho_{gh}, s)$  is the product of the local signs of  $L(E/K, \psi_1, s)$  and of  $L(E/K, \psi_2, s)$ , it follows that Hypothesis A in [DLR] holds.

Finally, the Heegner hypothesis also implies that there exists an integral ideal  $\mathfrak{N}$  in  $\mathcal{O}_K$ such that  $N_{K/\mathbb{Q}}(\mathfrak{N}) = N_f$ . We fix such an ideal for the remainder of the article.

Let X denote the modular curve associated in [LRdV, Theorem A] to the triple (E, K, c)and let

$$\pi_E:X\to E$$

be a modular parametrization of E.

Under our assumptions we let  $X(H)_c \subset X(H)$  denote the set of Heegner points attached to the order  $\mathcal{O}_c$ . After fixing a choice of  $P \in X(H)_c$ , any ring class character  $\psi$  of conductor c gives rise to a degree zero H-rational divisor on X by setting

(2) 
$$D_{\psi} := \begin{cases} \sum_{\sigma \in G} \psi^{-1}(\sigma) P^{\sigma}, & \text{if } \psi \neq 1, \\ \sum_{\sigma \in G} P^{\sigma} - \#G \cdot \infty & \text{if } \psi = 1, \end{cases}$$

where  $\infty$  is the cusp at infinity and  $G := \operatorname{Gal}(H/K)$ . The Heegner point attached to  $\psi$  is then defined to be

$$P_{\psi} := \pi_E(D_{\psi}) \in E(H) \otimes L.$$

Set  $\psi_0 := \psi_g/\psi_g'$  and note that this is again a ring class character of K. Fix an elliptic unit u in the field  $H_g$  cut out by  $\psi_0$  as in [DLR, §3.1] and define

$$u_{\psi_0} = \sum_{\sigma \in G} \psi_0^{-1}(\sigma) u^{\sigma} \in (\mathcal{O}_{H_g}^{\times}) \otimes L.$$

Let  $\omega_E$  denote the Néron differential 1-form of E. The pull-back of this form to X is a non-zero rational multiple of  $\omega_f$ , the canonical differential 1-form on X associated with f. Hence  $\pi_E^*(\omega_E) = c_E \cdot \omega_f$  for some  $c_E \in \mathbb{Q}^\times$ , usually referred to as the Manin constant of the modular parametrization. We refer to [ARS] for a survey of known results about the Manin constant.

Fix embeddings of  $H_q \subseteq H$  into  $\bar{\mathbb{Q}}_p \subset \mathbb{C}_p$  and let  $H_p$  denote the closure of H in  $\mathbb{C}_p$ . Let

$$\log_{E,p}: E(H_p) \longrightarrow H_p$$

denote the formal group logarithm associated with  $\omega_E$ . Let also

$$\log_p: \mathcal{O}_{H_p}^{\times} \longrightarrow \mathcal{O}_{H_p}$$

denote the usual p-adic logarithm. The main theorem of this note is the following.

**Theorem 1.1.** If either  $L'(E/K, \psi_1, 1)$  or  $L'(E/K, \psi_2, 1)$  vanishes, the iterated integrals attached to the triple (f, g, h) vanish identically.

Otherwise, there exists a finite extension L of  $\mathbb{Q}(\psi_q, \psi_h)$ , a scalar  $\lambda \in L^{\times}$  and test vectors

$$(\check{f},\check{\gamma}_{g_{\alpha}},\check{h})\in S_2(Np)_L[f]\times M_1(Np,\chi)^{\vee}_L[g_{\alpha}]\times M_1(Np,\chi)_L[h]$$

for which

$$\int_{\breve{\gamma}_{q_{\alpha}}} \breve{f} \cdot \breve{h} = \lambda \cdot \frac{\log_{E,p}(P_{\psi_1}) \log_{E,p}(P_{\psi_2})}{\log_p(u_{\psi_0})}.$$

The second main result of this note is a precise formula for the value of  $\lambda$  appearing in the above statement, in a particular setting.

**Theorem 1.2.** Assume  $N_f = D_K$  and  $g = h = \text{Eis}(1, \chi_K)$  are the Eisenstein series associated with the pair of Dirichlet characters  $(1, \chi_K)$ . Let  $h_K = |\text{Pic}(\mathcal{O}_K)|$  and  $g_K = |\text{Pic}(\mathcal{O}_K)|\text{Pic}^2(\mathcal{O}_K)|$  denote the class number and genus number of K. Then

$$\lambda = \frac{(p - a_p(f)\psi(\wp))^2}{p \cdot c_E^2} \cdot \frac{\lambda_0}{h_K g_K}$$

where:

$$\lambda_0 = \begin{cases} \frac{1}{p-1} & \text{if } \psi^2 = 1, \text{ that is to say, if } g \text{ is Eisenstein} \\ \frac{12}{p-(p+1)\psi^2(\bar{\wp})+\psi^4(\bar{\wp})} & \text{if } \psi^2 \neq 1, \text{ that is to say, if } g \text{ is cuspidal.} \end{cases}$$

The proof of these results is not a straight-forward generalization of [DLR] and [CR] respectively, mainly due to the following reasons:

- I. One of the main ingredients in the proof of Theorems 1.1 and 1.2 is the main result of Bertolini, Darmon and Prasanna in [BDP13], which is proved in loc. cit. under the assumption that p is a prime of good (although not necessarily ordinary) reduction.
- II. Another crucial ingredient in our proofs is the use of various p-adic L-functions, like Garrett-Hida's p-adic L-function constructed in [DR1], [DLR, Ch.2] and Bertolini-Darmon-Prasanna's p-adic L-function described in [BDP13, §5]. In the interpolation formulae relating these functions to critical values of classical L-series there appear certain auxiliary constants as for instance the ones denoted  $C_v$  in [DLR, Proposition 2.1] and also the one denoted  $\omega(f,\psi)$  in [BDP13, (5.1.11)]. When p is a prime of bad reduction, it is a rather subtle problem to analyze the p-adic valuation of these constants and its variation along Hida families of modular forms.
- III. In order to find an explicit formula for the scalar  $\lambda$  appearing in [DLR, Theorem 3.3], it was crucial in [CR] to assume that the levels of the three modular forms f, g and h were all equal, as otherwise the computations became too daunting. In the setting under study here, this assumption is not feasible because p divides the level of f but not the levels of g and h.

In order to overcome these difficulties we add an extra p-adic variable to the computations of [DLR, Ch. 2] by allowing the modular form f to vary in a Hida family. While in loc. cit. fwas fixed throughout the article, here we adopt the strategy of regarding it as a member of a Hida family **f** because the higher weight specialisations  $\mathbf{f}_k$ , k > 2, are (as opposite to f itself) old at p. This allows us to work essentially in a setting where p is a prime that behaves "as if it does not divide the level" and thus the necessary computations can then be performed more easily.

In the process of our proof, we prove in Theorem 6.1 a factorisation formula of two-variable p-adic L-functions involving Castellà's recent generalisation of [BDP13] that the reader may find of independent interest. Just as the analogous one-variable formula proved in DLR, Theorem 3.9 has found recently striking applications to the arithmetic of elliptic curves, we expect that our two-variable formula shall be also useful in this direction: cf. the forthcoming work of Bertolini and Darmon on Perrin-Riou's conjecture [PR93] for elliptic curves and Castellà's work in progress on the Iwasawa theory of elliptic curves over imaginary quadratic fields.

Acknowledgements. The two authors were supported by Spanish Grant MTM2015-63829-P. The first author acknowledges financial support by Spanish Ministry of Economy and Competitiveness. through the Severo Ochoa programme for Centres of Excellence in R&D (SEV-2015-0554). The second author acknowledges financial support from the Spanish Ministry of Economy and Competitiveness, through the María de Maeztu Programme for Units of Excellence in R&D (MDM-2014-0445). This project has received funding from the European Research Council (ERC) under the European Union's Horizon 2020 research and innovation programme (grant agreement No 682152). The authors also acknowledges financial support from the Spanish Ministry of Economy and Competitiveness, through the María de Maeztu Programme for Units of Excellence in R&D (MDM-2014-0445).

# 2. Garrett-Hida p-adic L-function

The starting point in the proof of Theorem 1.1 is the interpolation formula satisfied by the Garrett-Hida p-adic L-function associated with a triple of Hida families passing through f, gand h, that we now recall.

Fix a rational prime  $p \geq 3$  and set  $\Gamma = 1 + p\mathbb{Z}_p$  and let  $\Lambda = \mathbb{Z}_p[[\Gamma]]$  denote the Iwasawa algebra associated with  $\Gamma$ .

Define the formal scheme of weights as  $\mathcal{X} = \mathcal{X}_{\Lambda} := \mathrm{Spf}(\Lambda)$ . For any complete  $\mathbb{Z}_p$ -algebra R, the set of R-valued points of  $\mathcal{X}$  may be identified with

$$\mathcal{X}(R) = \operatorname{Hom}_{\operatorname{cont}}(\Lambda, R) \simeq \operatorname{Hom}(\Gamma, R^{\times}).$$

As usual, the set  $\mathbb{Z}_{\geq 2}$  can be embedded in  $\mathcal{X}(\mathbb{Z}_p)$  by identifying an integer k with the character  $\nu_k : \Gamma \longrightarrow \mathbb{Z}_p^{\times}$  sending z to  $z^k$ . We let  $\mathcal{X}^{\text{cl}}$  denote the set of such classical points.

Given a finite flat extension  $\tilde{\Lambda}$  of  $\Lambda$ , there is a natural projection  $\mathcal{X}_{\tilde{\Lambda}} \to \mathcal{X}$  induced by the inclusion  $\Lambda \hookrightarrow \tilde{\Lambda}$  and we define  $\mathcal{X}_{\tilde{\Lambda}}^{\text{cl}}$  as the pull-back of  $\mathcal{X}^{\text{cl}}$  under this map. The set  $\mathcal{X}_{\tilde{\Lambda}}^{\text{cl}}$  is called set of *arithmetic primes* of  $\tilde{\Lambda}$ . By definition, any  $\nu \in \mathcal{X}_{\tilde{\Lambda}}^{\text{cl}}$  lies above  $\nu_k \in \mathcal{X}^{\text{cl}}$  for some  $k = k(\nu) \ge 2$ , which we refer to as the weight of  $\nu$ .

**Definition 2.1.** Let  $N_f \geq 1$  be an arbitrary integer not divisible by p and  $\chi$  be a Dirichlet character modulo  $N_f$ . A Hida family of tame level  $N_f$  and tame character  $\chi$  is a quadruple  $(\Lambda_f, \mathcal{U}_f, \mathcal{U}_f^{\mathrm{cl}}, \mathbf{f})$  where

- $\Lambda_f$  is a finite flat extension of  $\Lambda$ ,  $\mathcal{U}_f \subset \mathcal{X}_{\Lambda_f}$  is a rigid-analytic open subvariety,
- $\mathcal{U}_f^{\mathrm{cl}} \subset \mathcal{U}_f(\mathbb{C}_p)$  is a dense subset for the rigid analytic topology, and

•  $\mathbf{f} = \sum \mathbf{a}_n q^n \in \Lambda_f[[q]]$  is a formal q-series such that for all  $\nu \in \mathcal{U}_f^{\mathrm{cl}}$  of weight  $k = k(\nu)$ ,

$$\nu(\mathbf{f}) = \sum_{n \ge 1} \nu(\mathbf{a}_n) q^n \in S_k(pN_f, \chi \omega^{k_0 - k})$$

is the q-expansion of an ordinary eigenform of weight k, level  $pN_f$  and nebentype  $\chi \omega^{k_0-k}$ , for some  $k_0 \in \{0, \ldots, p-2\}$ .

Here  $\omega: (\mathbb{Z}/p\mathbb{Z})^{\times} \to \mu_{p-1}$  stands for the Teichmüller character of conductor p.

Remark 2.2. In the literature one often allows a richer plethora of classical specializations, including twists by characters of p-power conductor of  $\Gamma$ . For our purposes it suffices to restrict to the case in which this character is trivial. Given an arithmetic prime  $\nu \in \mathcal{U}_f^{\text{cl}}$  over the classical point  $\nu_k \in \mathcal{X}$ , it will often be harmless to abuse notation and write  $\mathbf{f}_k$  instead of  $\nu(\mathbf{f})$ .

Remark 2.3. In order to simplify further the notations, we shall always restrict to those  $k \in \mathbb{Z}_{\geq 1}$  for which  $k \equiv k_0 \pmod{p-1}$ , so that the nebentype character of  $\mathbf{f}_k$  will be exactly  $\chi$ . Under this condition, the specialisation  $\mathbf{f}_k$  of the Hida family  $\mathbf{f}$  at a point of weight k > 2 is a classical eigenform which is always old at p (cf. e.g. [Ho07, Lemma 2.1.5]) and there exists an eigenform  $f_k$  of level N such that  $\mathbf{f}_k(q) = f_k(q) - \beta_{f_k} f_k(q^p)$  is the ordinary p-stabilization of  $f_k$ . Here  $\beta_{f_k}$  is the single non-unit root of the Hecke polynomial  $T^2 - a_p(f_k)T + \chi(p)p^{k-1}$ . The Hecke operator  $U_p$  acts on  $\mathbf{f}_k$  with eigenvalue  $\alpha_{f_k}$ , the single p-adic unit root of the above polynomial.

If k=2, then  $\mathbf{f}_2$  is a classical eigenform that may either be old or new p. In the former case there exists as above an eigenform  $f_2$  of level N such that  $\mathbf{f}_2$  is the ordinary p-stabilisation of  $f_2$ . In the latter case we simply set  $f_2 = \mathbf{f}_2$  and  $\alpha_{f_2} = a_p(f_2)$ ,  $\beta_{f_2} = 0$ .

If k = 1, the p-adic modular form  $\mathbf{f}_1$  may be classical or not: see [DG] for a discussion of this phenomenon.

Recall now the three classical eigenforms

$$f \in S_2(pN_f), \quad g \in M_1(N_g, \bar{\chi}), \quad h \in M_1(N_h, \chi)$$

introduced at the beginning of this note. Choose eigenforms

$$\check{f} \in S_2(Np)[f], \quad \check{g} \in M_1(N,\bar{\chi})[g], \quad \check{h} \in M_1(N,\chi)_L[h]$$

with respect to the good Hecke operators, having the same eigenvalues as f, g and h respectively.

Since g and h have weight 1, the roots of the Hecke polynomials

$$T^2 - a_p(g)T + \bar{\chi}(p) = (T - \alpha_g)(T - \beta_g), \quad T^2 - a_p(h)T + \chi(p) = (T - \alpha_h)(T - \beta_h)$$

are all roots of unity. Fix arbitrary orderings  $(\alpha_a, \beta_a)$ ,  $(\alpha_b, \beta_b)$  of these roots and set

$$g_{\alpha}(q) := g(q) - \beta_g g(q^p), \quad h_{\alpha}(q) := h(q) - \beta_h h(q^p),$$

$$\check{g}_{\alpha}(q) := \check{g}(q) - \beta_g \check{g}(q^p), \quad \check{h}_{\alpha}(q) := \check{h}(q) - \beta_h \check{h}(q^p).$$

Theorems of Hida [Hi86] and Wiles [Wi88] ensure that there exist Hida families  $\check{\mathbf{f}}, \check{\mathbf{g}}, \check{\mathbf{h}}$  of tame level N such that, for suitable arithmetic primes of weights 2, 1 and 1 we have

$$\check{\mathbf{f}}_2(q) = \check{f}(q), \quad \check{\mathbf{g}}_1(q) = \check{g}_{\alpha}(q), \quad \check{\mathbf{h}}_1(q) = \check{h}_{\alpha}(q).$$

Define a dense subset of  $\mathcal{U}_f \times \mathcal{U}_q \times \mathcal{U}_h$  as follows:

$$\mathcal{U}_{fgh}^{\text{cl}} := \{(k,\ell,m) \in \mathcal{U}_f^{\text{cl}} \times \mathcal{U}_g^{\text{cl}} \times \mathcal{U}_h^{\text{cl}} \mid \ell \geq k+m \}.$$

In [DR1], the authors constructed a three-variable p-adic L-function on  $\mathcal{U}_f \times \mathcal{U}_g \times \mathcal{U}_h$  which is determined by an interpolation formula relating the values at points in  $\mathcal{U}_{fgh}^{\text{cl}}$  to central critical values of classical L-series: we refer to [DLR, prop. 2.3] for full-fledged details.

For our purposes it suffices to restrict the above p-adic L-function to  $\mathcal{U}_{fg} := \mathcal{U}_f \times \mathcal{U}_g$ , regarded as a rigid-analytic sub-variety of  $\mathcal{U}_f \times \mathcal{U}_g \times \mathcal{U}_h$  by setting the third variable to be the fixed point of weight one giving rise to  $\check{h}_{\alpha}$ . Set also  $\mathcal{U}_{fg}^{\text{cl}} := \{(k,\ell) \in \mathcal{U}_f^{\text{cl}} \times \mathcal{U}_g^{\text{cl}}, \ell \geq k+1\}$ . With these notations, the constructions of [DR1] and [DLR, Ch. 2] give rise to a two-variable

p-adic L-function

$$\mathscr{L}_p{}^g(\check{\mathbf{f}},\check{\mathbf{g}},\check{h}):\mathcal{U}_{fg}\longrightarrow\mathbb{C}_p$$

satisfying the following interpolation property for all  $(k,\ell) \in \mathcal{U}_{fg}^{\mathrm{cl}}$ :

$$\mathcal{L}_{p}^{g}(\check{\mathbf{f}},\check{\mathbf{g}},\check{h})\cdot\mathcal{L}_{p}^{g}(\check{\mathbf{f}}^{*},\check{\mathbf{g}}^{*},\check{h}^{*})(k,\ell) = \mathfrak{e}_{\mathrm{GH}}(k,\ell)^{2}\cdot\mathfrak{a}_{\mathrm{GH}}(k,\ell)\cdot\mathfrak{f}_{\mathrm{GH}}(k,\ell)\cdot$$

$$\times \prod_{v|N\infty} C_{v}(\check{f}_{k},\check{g}_{\ell},\check{h})\cdot\frac{L(f_{k}\otimes g_{\ell}\otimes h,(k+\ell-1)/2)}{\langle g_{\ell}^{*},g_{\ell}^{*}\rangle^{2}}$$
(3)

where:

- For any modular form or Hida family  $\phi$  of character  $\chi$ , we set  $\phi^* = \phi \otimes \bar{\chi}$ ;
- $L(f_k, \otimes g_\ell \otimes h, s)$  is Garrett's triple-product complex L-function;
- $\mathfrak{a}_{GH}(k,\ell) = (\frac{\ell+k-3}{2})!^2(\frac{\ell-k-1}{2})!^2 \cdot \pi^{-2\ell};$
- $\mathfrak{e}_{GH}(k,\ell) = \mathcal{E}(g_{\ell}, f_k, h)/\mathcal{E}_0(g_{\ell})\mathcal{E}_1(g_{\ell}) = \mathcal{E}(g_{\ell}^*, f_k^*, h^*)/\mathcal{E}_0(g_{\ell}^*)\mathcal{E}_1(g_{\ell}^*)$  is the Euler factor introduced in [DR1, Theorem 1.3];
- $\mathbf{f}_{GH}(k,\ell) = A \cdot 2^B$  where  $A = A(N) \in \mathbb{Q}^{\times}$  and  $B = B(k,\ell)$  is linear in  $k,\ell$ ;
- $C_v(\check{f}_k, \check{g}_\ell, \check{h})$  is a constant in the field generated by the fourier coefficients of  $\check{f}_k, \check{g}_\ell$  and  $\check{h}$ ; thanks to the work of Watson and Woodbury [Wo1], when the three levels of the modular forms are square-free there exists a choice of test vectors  $(\hat{f}_k, \check{g}_\ell, \check{h})$  such that the constants

$$(4) W_v := C_v(\check{f}_k, \check{g}_\ell, \check{h})$$

satisfy that  $W_v$  lies in  $\mathbb{Q}^{\times}$  for all  $v \mid N$  and only depends on N and v but not on the weights  $(k, \ell)$  (see [DLR, remark 2.2 and remark 2.4]).

When we evaluate this function at the point  $(2,1) \in \mathcal{U}_{fg}$ , which lies outside the region of classical interpolation, [DLR, proposition 2.6] asserts that:

(5) 
$$\mathscr{L}_{p}^{g}(\breve{\mathbf{f}},\breve{\mathbf{g}},\breve{h})(2,1) = \int_{\breve{\gamma}_{\alpha}} \breve{f}_{2} \cdot \breve{h}$$

for a suitable choice of linear functional  $\check{\gamma}_{\alpha}: S_1(Np,\chi)_L[\check{g}_{\alpha}^*] \to L$ , with  $L = \mathbb{Q}(f_2,g_1,h)$ .

# 3. Castellà's two-variable p-adic L-function

Consider the Hida family  $\mathbf{f}$  as in the previous section. For any weight  $\kappa \in \mathcal{U}_f^{\text{cl}}$  such that  $\kappa \equiv 2 \pmod{p-1}$  let  $f_{\kappa}$  denote the modular form of weight  $\kappa$  whose p-stabilization equals the weight  $\kappa$  specialization of  $\mathbf{f}$ , as in Remark 2.3.

Let  $K = \mathbb{Q}(\sqrt{-D_K})$ ,  $D_K \geq 7$ , be an imaginary quadratic field. We keep the same assumptions as in the introduction, so in particular we assume the Heegner hypothesis and we let  $p = \wp \bar{\wp}$  be an odd prime that splits in K.

A Hecke character  $\psi$  of K of infinity type  $(\kappa_1, \kappa_2) \in \mathbb{Z} \times \mathbb{Z}$  is a continuous homomorphism

$$\psi: \mathbb{A}_K^{\times} \longrightarrow \mathbb{C}^{\times}$$

satisfying  $\psi(\alpha \cdot x \cdot z_{\infty}) = \psi(x) z_{\infty}^{\kappa_1} \bar{z}_{\infty}^{\kappa_2}$  for every  $\alpha \in K^{\times}$ ,  $x \in \mathbb{A}_K^{\times}$  and  $z_{\infty} \in K_{\infty}^{\times} = \mathbb{C}^{\times}$ . With this convention, the Hecke character  $\mathbf{N}_K$  induced by the norm from K to  $\mathbb{Q}$  has infinity type (1,1). A Hecke character  $\psi$  of infinity type  $(\kappa_1, \kappa_2)$  is called *anticyclotomic* if its central character satisfies  $\psi_{|\mathbb{A}_{\mathbb{Q}}^{\times}} = \mathbf{N}_{K}^{\kappa_{1}+\kappa_{2}}$ .

Let  $\mathfrak{c} \subseteq \mathcal{O}_K$  be an integral ideal and set  $N = \operatorname{lcm}(N_f, D_K N_{K/\mathbb{Q}}(\mathfrak{c}))$ . Let  $\Sigma_{\operatorname{cc}}$  denote the set of anticyclotomic Hecke characters of K of conductor dividing N. For any Hecke character  $\psi \in \Sigma_{cc}$  of infinity type  $(\kappa_1, \kappa_2)$  such that  $\kappa = \kappa_1 + \kappa_2 \in \mathcal{U}_f^{cl}$ , let  $L(f_\kappa, \psi, s)$  denote the Lfunction associated with the twist of  $f_{\kappa}$  by  $\psi$ , i.e. the L-function attached to the compatible system of Galois representations afforded by the tensor product  $\varrho_{f_{\kappa}|G_K} \otimes \psi$  of the (restriction to  $G_K$  of) the Galois representations attached to  $f_{\kappa}$  and the character  $\psi$ .

As usual,  $L(f_{\kappa}, \psi, s) = \prod_{q} L^{(q)}(q^{-s})$  is defined as a product of Euler factors ranging over the set of prime numbers. The Euler factors at the primes q such that  $q \nmid N$  are the same as that of the Rankin L-series  $L(\theta_{\psi} \otimes f_{\kappa}, s)$ , but may differ at the primes q such that  $q \mid N$ . Since  $\psi$  is anticyclotomic, s=0 is the central critical point of the functional equation satisfied by  $L(f_{\kappa}, \psi^{-1}, s).$ 

For a given  $\kappa \in \mathcal{U}_f^{\text{cl}}$ , let  $\Sigma_{\kappa,\text{cc}} \subset \Sigma_{\text{cc}}$  be the subset of anticyclotomic Hecke characters in  $\Sigma_{\text{cc}}$ of infinity type  $(\kappa_1, \kappa_2)$  with  $\kappa_1 + \kappa_2 = \kappa$ . This set is naturally the disjoint union of the three subsets

$$\Sigma_{\kappa, cc}^{(1)} = \{ \psi \in \Sigma_{\kappa, cc} \text{ of infinity type } (1, 1) \},$$

$$\Sigma_{\kappa,cc}^{(2)} = \{ \psi \in \Sigma_{\kappa,cc} \text{ of infinity type } (\kappa + \lambda, -\lambda), \lambda \in \mathbb{Z}_{\geq 0} \}$$

and

$$\Sigma_{\kappa, cc}^{(2')} = \{ \psi \in \Sigma_{\kappa, cc} \text{ of infinity type } (-\lambda, \lambda + \kappa), \lambda \in \mathbb{Z}_{\geq 0} \}.$$

Each of the three sets  $\Sigma_{\kappa,cc}^{(1)}$ ,  $\Sigma_{\kappa,cc}^{(2)}$  and  $\Sigma_{\kappa,cc}^{(2')}$  is dense in the completion  $\hat{\Sigma}_{\kappa,cc}$  of  $\Sigma_{\kappa,cc}$  with respect to the *p*-adic compact open topology as explained in [BDP13, §5.2]. In loc. cit. the authors construct a p-adic L-function on  $\tilde{\Sigma}_{\kappa,cc}$ , denoted  $\mathscr{L}_p(f_{\kappa}/K)$ , interpolating the algebraic parts of the central critical values of  $L(f_{\kappa}, \psi^{-1}, s)$  at characters  $\psi$  in  $\Sigma_{\kappa, cc}^{(2)}$ .

In [Ca1, §1], Castellà constructs a two-variable p-adic L-function  $\mathcal{L}_p(\mathbf{f}/K)(\kappa,\psi)$  on the completion  $\hat{\Sigma}_{cc}$  of  $\Sigma_{cc}$ , lifting the one-variable p-adic L-function of Bertolini-Darmon-Prasanna. To state this more precisely, note that  $\hat{\Sigma}_{cc}$  is naturally fibered over  $\mathcal{U}_f$  and the fiber at a given point  $\kappa \in \mathcal{U}_f^{\text{cl}}$  is  $\hat{\Sigma}_{\kappa,\text{cc}}$ . Then for any  $\kappa \in \mathcal{U}_f^{\text{cl}}$  we have

$$\mathscr{L}_p(\mathbf{f}/K)(\kappa,-)^2 = \mathscr{L}_p(f_{\kappa}/K).$$

Castellà's p-adic L-function is characterized by the following interpolation property: let  $\kappa \in \mathcal{U}_f$  and let  $\psi \in \Sigma_{\kappa, cc}^{(2)}$  be a character of infinity type  $(\kappa + \lambda, -\lambda)$ . Then [Ca1, Theorem 1.4] asserts that

(6) 
$$\mathscr{L}_p(\mathbf{f}/K)(\kappa,\psi)^2 = \mathfrak{a}_{\mathrm{BDP}}(\psi) \cdot \mathfrak{e}_{\mathrm{BDP}}(\kappa,\psi)^2 \cdot \mathfrak{f}_{\mathrm{BDP}}(\kappa,\psi) \cdot \frac{\Omega_p^{2\kappa+4\lambda}}{\Omega^{2\kappa+4\lambda}} \cdot L(f_\kappa,\psi^{-1},0),$$

where

- $\mathfrak{a}_{BDP}(\psi) = \lambda!(\kappa + \lambda 1)!\pi^{\kappa + 2\lambda 1}$
- $\mathfrak{e}_{\mathrm{BDP}}(\kappa, \psi) = (1 \alpha_{f_{\kappa}} \psi^{-1}(\bar{\wp}))(1 \beta_{f_{\kappa}} \psi^{-1}(\bar{\wp})),$   $\mathfrak{f}_{\mathrm{BDP}}(\kappa, \psi) = \left(\frac{2}{c\sqrt{D_{K}}}\right)^{\kappa + 2\lambda 1} \cdot 2^{\sharp\{q|(D_{K}, N_{E})\}} \cdot \omega(f_{\kappa}, \psi)^{-1}$

with  $\omega(f_{\kappa}, \psi)$  as defined in [BDP13, (5.1.11)].

If  $\psi$  is a finite order anticyclotomic character of conductor  $c \mid \mathfrak{c}$ , then  $\psi^{-1} \mathbf{N}_K$  is a character of infinity type (1,1) and hence lies outside the region of interpolation. The main theorems of [BDP13] and [Ca2, Theorem 2.10] assert that for such  $\psi$  the following formula holds:

(7) 
$$\mathscr{L}_p(\mathbf{f}/K)(2,\psi\mathbf{N}_K) = \pm \mathfrak{f}_p(f,\psi) \cdot \log_{\omega_f}(D_\psi)$$

where  $f_p(f, \psi) = (1 - \alpha_{f_2} \psi^{-1}(\bar{\wp}) p^{-1}) (1 - \beta_{f_2} \psi^{-1}(\bar{\wp}) p^{-1})$  and  $D_{\psi}$  is the class in Jac(X) of the divisor introduced in (2).

Since  $\pi_E^*(\omega_E) = c_E \cdot \omega_f$  and  $P_{\psi} = \pi_E(D_{\psi})$ , it follows that

$$\log_{\omega_f}(D_{\psi}) = \frac{1}{c_E} \log_{\pi_E^*(\omega_E)}(D_{\psi}) = \frac{1}{c_E} \log_{E,p}(P_{\psi}),$$

and hence

(8) 
$$\mathscr{L}_p(\mathbf{f}/K)(2, \psi \mathbf{N}_K) = \pm \mathfrak{f}_p(f, \psi) \cdot \frac{1}{c_E} \log_{E,p}(P_{\psi}).$$

## 4. Katz p-adic L-function

Keep the notations and assumptions of the previous section. Denote by  $\Sigma_K$  the set of Hecke characters of K of conductor dividing  $\mathfrak{c}$  and define

$$\Sigma_K^{(1)} = \{ \psi \in \Sigma \text{ of infinity type } (\kappa_1, \kappa_2), \kappa_1 \leq 0, \kappa_2 \geq 1 \},$$

$$\Sigma_K^{(2)} = \{ \psi \in \Sigma \text{ of infinity type } (\kappa_1, \kappa_2), \kappa_1 \ge 1, \kappa_2 \le 0 \}.$$

For all  $\psi \in \Sigma_K^{(2)}$ , s = 0 is a critical point for the Hecke *L*-function  $L(\psi^{-1}, s)$ , and Katz's *p*-adic *L*-function is constructed by interpolating the (suitably normalized) values  $L(\psi^{-1}, 0)$ as  $\psi$  ranges over  $\Sigma_K^{(2)}$ .

More precisely, let  $\hat{\Sigma}_K$  denote the completion of  $\Sigma_K^{(2)}$  with respect to the compact open topology on the space of functions on a certain subset of  $\mathbf{A}_{K}^{\times}$ , as described in [BDP13, §5.2]. By the work of Katz [Ka76], there exists a p-adic analytic function

$$\mathscr{L}_p(K): \hat{\Sigma}_K \longrightarrow \mathbb{C}_p$$

which is characterized by the following interpolation property: for all  $\psi \in \Sigma_K^{(2)}$  of infinity type  $(\kappa_1, \kappa_2),$ 

(9) 
$$\mathscr{L}_p(K)(\psi) = \mathfrak{a}_K(\psi) \cdot \mathfrak{e}_K(\psi) \cdot \mathfrak{f}_K(\psi) \frac{\Omega_p^{\kappa_1 - \kappa_2}}{\Omega^{\kappa_1 - \kappa_2}} L_{\mathfrak{c}}(\psi^{-1}, 0)$$

where

- $L_{\mathfrak{c}}(\psi^{-1},s)$  is Hecke's L-function associated with  $\psi^{-1}$  with the Euler factors at primes dividing  $\mathfrak{c}$  removed,
- $\Omega_p \in \mathbb{C}_p^{\times}$  is a p-adic period attached to K, as defined in [BDP13, (140)],
- $\Omega \in \mathbb{C}^{\times}$  is the complex period associated with K as defined in [BDP13, (137)],
- $\bullet \ \mathfrak{a}_{\mathrm{K}}(\psi) = \frac{(\kappa_{1}-1)!}{\pi^{\kappa_{2}}}$   $\bullet \ \mathfrak{e}_{\mathrm{K}}(\psi) = (1 \frac{\psi(\wp)}{p})(1 \psi^{-1}(\bar{\wp}))$   $\bullet \ \mathfrak{f}_{\mathrm{K}}(\psi) = \left(\frac{\sqrt{D_{\mathrm{K}}}}{2}\right)^{\kappa_{2}}.$

The following result is commonly known as Katz's Kronecker p-adic limit formula. It computes the value of  $\mathcal{L}_p(K)$  at a finite order character  $\psi$  of  $G_K$ , which lies outside the region of interpolation (cf. [Ka76, §10.4, 10.5]):

(10) 
$$L_p(K)(\psi) = \mathfrak{f}_p(\psi) \cdot \log_p(u_\psi),$$

where

(11) 
$$\mathfrak{f}_p(\psi) = \begin{cases} \frac{1}{2} (\frac{1}{p} - 1) & \text{if } \psi = 1\\ \frac{-1}{24c} (1 - \psi(\bar{\wp})) (1 - \frac{\psi(\bar{\wp})}{p}) & \text{if } \psi \neq 1. \end{cases}$$

Here c > 0 is the smallest positive integer in the conductor ideal of  $\psi$ .

## 5. Proof of Theorem 1.1

Let

$$f \in S_2(pN_f), \qquad g = \theta_{\psi_g} \in M_1(D_K N(\mathfrak{c}_g), \overline{\chi}), \qquad h = \theta_{\psi_h} \in M_1(D_K N(\mathfrak{c}_h), \chi)$$

be the three modular forms considered in the introduction, and we keep the assumptions as therein. Remember that  $p = \wp \overline{\wp}$  is a prime that splits in  $\mathcal{O}_K$  and does not divide  $cN_f$ .

Let **f** and **g** be cuspidal Hida families passing through f and  $g_{\alpha}$ , that is to say, satisfying  $\mathbf{f}_2 = f$  and  $\mathbf{g}_1 = g_{\alpha}$ . The Hida family  $\mathbf{f}$  is unique thanks to the classical work of Hida [Hi86]. The Hida family g is also unique thanks to the classicality hypothesis assumed in the introduction, by a recent result of Bellaiche and Dimitrov [BeDi]. This family is in fact CM and its higher weight classical specializations admit an explicit description that we now review.

Take a Hecke character  $\lambda$  with image in  $\mathbb{Z}_p^{\times}$  of infinity type (0,1) and conductor  $\bar{\wp}$ . Following [Hi93, p. 235-236], with a slight modification, we define the Hecke characters  $\psi_{g,\ell-1}$ , of conductor  $\mathfrak{c}_q$ , as follows:

$$\psi_{g,\ell-1}(\mathfrak{q}) := \psi_g(\mathfrak{q}) \langle \lambda \rangle^{\ell-1}(\mathfrak{q}), \qquad \text{for } \mathfrak{q} \neq \bar{\wp}$$

$$\psi_{g,\ell-1}(\bar{\wp}) := \overline{\chi}(p) \frac{p^{\ell-1}}{\psi_{g,\ell-1}(\wp)}.$$

For every  $\ell \in \mathcal{U}_q^{\text{cl}}$  one then has that  $\mathbf{g}_{\ell}$  is the ordinary *p*-stabilisation of  $g_{\ell} := \theta_{\psi_{q,\ell-1}}$ . Using the notations introduced above, note that we have

$$\alpha_{q_{\ell}} = \psi_{q,\ell-1}(\wp), \qquad \beta_{q_{\ell}} = \psi_{q,\ell-1}(\bar{\wp}).$$

Recall the three ring class characters  $\psi_0$ ,  $\psi_1$ ,  $\psi_2$  defined in the introduction. Recall also that  $N_K$  stands for the Hecke character of infinity type (1,1) induced by the norm from K to Q. Associated with them there are three additional families of characters which play a key role in the proof:

(12) 
$$\Phi_q(\ell) := (\psi_{q\ell-1}^2 \chi)^{-1} \mathbf{N}_K^{\ell}, \qquad \text{for which } \Phi_q(1) = \psi_0' \mathbf{N}_K$$

(12) 
$$\Phi_{g}(\ell) := (\psi_{g,\ell-1}^{2}\chi)^{-1}\mathbf{N}_{K}^{\ell}, \qquad \text{for which } \Phi_{g}(1) = \psi_{0}^{\prime}\mathbf{N}_{K}$$
(13) 
$$\Psi_{gh}(k,\ell) := (\psi_{g,\ell-1}\psi_{h})^{-1}\mathbf{N}_{K}^{\frac{k+\ell-1}{2}}, \qquad \text{for which } \Psi_{gh}(2,1) = \psi_{1}^{\prime}\mathbf{N}_{K}$$
(14) 
$$\Psi_{gh^{\prime}}(k,\ell) := (\psi_{g,\ell-1}\psi_{h}^{\prime})^{-1}\mathbf{N}_{K}^{\frac{k+\ell-1}{2}}, \qquad \text{for which } \Psi_{gh}(2,1) = \psi_{2}^{\prime}\mathbf{N}_{K}$$

(14) 
$$\Psi_{gh'}(k,\ell) := (\psi_{g,\ell-1}\psi'_h)^{-1}\mathbf{N}_K^{\frac{k+\ell-1}{2}}, \quad \text{for which } \Psi_{gh}(2,1) = \psi'_2\mathbf{N}_K.$$

Notice that  $\Phi_g(\ell)$  is of infinity type  $(\kappa_1, \kappa_2) = (\ell, -\ell + 2)$ , while  $\Psi_{gh}(k, \ell)$  and  $\Psi_{gh'}(k, \ell)$  are of infinity type  $(\kappa + \lambda, -\lambda)$  with  $\kappa = k$  and  $\lambda = (\ell - k - 1)/2 \ge 0$  because  $\ell \ge k + 1$ .

**Definition 5.1.** Let L be a number field. A function

$$\mathfrak{f}:\mathcal{U}_g^{\mathrm{cl}}\to\overline{\mathbb{Q}}$$

is L-admissible if it extends to a meromorphic Iwasawa function on  $\mathcal{U}_g$  having no pole at 1 and satisfying  $\mathfrak{f}(1) \in L^{\times}$ . Similarly, a function

$$\mathfrak{f}:\mathcal{U}^{\mathrm{cl}}_{fg} o\overline{\mathbb{Q}}$$

is L-admissible if it extends to a meromorphic Iwasawa function on  $\mathcal{U}_{fq}$  having no pole at (2,1) and satisfying  $\mathfrak{f}(2,1) \in L^{\times}$ .

Set  $\mathfrak{a}_{\mathrm{Pet}}(\ell) = \frac{(\ell-1)!}{\pi^{\ell}}$ . By [DLR, Lemma 3.7 and (53)] there exists a  $\mathbb{Q}(\psi_g, \psi_h)$ -admissible function  $\mathfrak{f}_{\mathrm{Pet}}$  such that

(15) 
$$\langle g_{\ell}^*, g_{\ell}^* \rangle = \mathfrak{a}_{\mathrm{Pet}}(\ell) \cdot \mathfrak{f}_{\mathrm{Pet}}(\ell) \cdot L(\Phi_g(\ell)^{-1}, 0).$$

**Lemma 5.2.** The following formulae hold true:

$$\begin{split} \Phi_g(\ell)(\wp) &= \frac{\beta_{g_\ell}^2 \chi(p)}{p^{\ell-2}}, \qquad \qquad \Phi_g(\ell)(\bar{\wp}) = \frac{p^\ell}{\beta_{g_\ell}^2 \chi(p)}, \\ \Psi_{gh}(k,\ell)(\bar{\wp}) &= \frac{p^{\frac{\ell+k-1}{2}}}{\beta_{g_\ell}\beta_h}, \qquad \qquad \Psi_{gh'}(k,\ell)(\bar{\wp}) = \frac{p^{\frac{\ell+k-1}{2}}}{\beta_{g_\ell}\alpha_h}. \end{split}$$

*Proof.* This follows from a straight-forward computation.

**Lemma 5.3.** The following equalities hold true:

$$\begin{split} & \mathbf{e}_K(\Phi_g(\ell))^2 \cdot \mathbf{e}_{\mathrm{G}H}(k,\ell)^2 = \mathbf{e}_{\mathrm{BDP}}(k,\Psi_{gh}(k,\ell))^2 \mathbf{e}_{\mathrm{BDP}}(k,\Psi_{gh'}(k,\ell))^2, \\ & \mathbf{a}_{\mathrm{G}H}(k,\ell) \cdot \mathbf{a}_{\mathrm{K}}(\Phi_g(\ell))^2 = \mathbf{a}_{\mathrm{B}DP}(\Psi_{gh}(k,\ell)) \cdot \mathbf{a}_{\mathrm{B}DP}(\Psi_{gh'}(k,\ell)) \cdot \mathbf{a}_{\mathrm{Pet}}(\ell)^2. \end{split}$$

*Proof.* This again follows from a straight-forward computation using Lemma 5.2 for the first equality.  $\Box$ 

Recall the factor  $\omega(f, \Psi)$  appearing in (6) and defined in [BDP13, equation (5.1.11)]. Recall also from the remarks around the Heegner hypothesis in the introduction the integral ideal  $\mathfrak{N}$  of  $\mathcal{O}_K$  such that  $N_{K/\mathbb{Q}}(\mathfrak{N}) = N_f$ .

Lemma 5.4. The following identities hold true:

$$\omega(f, \Psi_{gh}(k, \ell)) = (-1)^{\frac{\ell - 1}{2}} \frac{\psi_{g, \ell - 1}(\mathfrak{N})\psi_h(\mathfrak{N})}{N_f^{\frac{\ell - 1}{2}}}, \qquad \omega(f, \Psi_{gh}(k, \ell)) = (-1)^{\frac{\ell - 1}{2}} \frac{\psi_{g, \ell - 1}(\mathfrak{N})\psi_h'(\mathfrak{N})}{N_f^{\frac{\ell - 1}{2}}}.$$

In particular we have

$$\omega(f, \Psi_{ab}(2,1)) = \psi_1(\mathfrak{N})$$
 and  $\omega(f, \Psi_{ab}(2,1)) = \psi_2(\mathfrak{N}).$ 

*Proof.* We limit ourselves to compute  $\omega(f, \Psi_{gh}(k, \ell))$ , as the other case is very similar. For simplicity, call  $\Psi := \Psi_{gh}(k, \ell)$  the Hecke character of infinity type  $(\kappa + \lambda, -\lambda)$ , where  $\kappa = k$  and  $\lambda = (\ell - k - 1)/2$ . Define

$$\Psi_{\lambda} := \Psi \mathbf{N}_{K}^{\lambda}.$$

Choose an ideal  $\mathfrak{b} \subset \mathcal{O}_c$  relatively prime to  $N_f c$  and an element  $b_N \in \mathcal{O}_c$  such that  $\mathfrak{b}\mathfrak{N} = (b_N)$ . Since  $\omega_f = 1$  and  $\varepsilon_f = 1$  in our case, [BDP13, equation (5.1.11)] shows that

(16) 
$$\omega(f, \Psi) = \Psi_{\lambda}(\mathfrak{b}) \cdot (-1)^{\kappa/2 + \lambda} N^{\kappa/2 + \lambda} b_N^{-\kappa - 2\lambda}.$$

By [BDP13, equation (5.1.2)] it then follows that

$$\Psi_{\lambda}(\mathfrak{b})\Psi_{\lambda}(\mathfrak{N}) = \Psi_{\lambda}((b_N)) = b_N^{\kappa + 2\lambda}$$

Substituting into (16) we find:

$$\omega(f, \Psi) = (-1)^{\kappa/2 + \lambda} \frac{N_f^{\kappa/2 + \lambda}}{\Psi_{\lambda}(\mathfrak{N})} = (-1)^{\kappa/2 + \lambda} \frac{N_f^{\kappa/2}}{\Psi(\mathfrak{N})}.$$

Using now the definition of  $\Psi_{qh}(k,\ell)$  and the fact that  $\kappa + 2\lambda = \ell - 1$  we find

$$\omega(f, \Psi) = (-1)^{\frac{\ell-1}{2}} \cdot \frac{\psi_{g,\ell-1}(\mathfrak{N})\psi_h(\mathfrak{N})N_f^{k/2}}{N_f^{\frac{\ell+k-1}{2}}} = (-1)^{\frac{\ell-1}{2}} \cdot \frac{\psi_{g,\ell-1}(\mathfrak{N})\psi_h(\mathfrak{N})}{N_f^{\frac{\ell-1}{2}}},$$

as claimed.  $\Box$ 

Let

$$L_0 = K(\psi_g, \psi_h, \tau(\chi), \sqrt{N_g}, \sqrt{N_h})$$

denote the extension of  $\mathbb Q$  generated by the values of the Hecke characters  $\psi_g$  and  $\psi_h$ , the Gauss sum associated with the Dirichlet character  $\chi$  and the square-roots of  $N_g$  and  $N_h$ .

Fix test vectors  $\check{\mathbf{f}}$ ,  $\check{\mathbf{g}}$  and  $\check{h}$  as in §2, giving rise to the p-adic L-function  $\mathcal{L}_p{}^g(\check{\mathbf{f}}, \check{\mathbf{g}}, \check{h})$  satisfying the good properties discussed around (4).

**Theorem 5.5.** There exist a quadratic extension  $L/L_0$  and a L-admissible function  $\mathfrak{f}$  on  $\mathcal{U}_{fg}$  such that the following factorization of two-variable p-adic L-functions holds:

$$\begin{aligned} \mathscr{L}_{p}^{g}(\check{\mathbf{f}}, \check{\mathbf{g}}, \check{h})(k, \ell) \cdot \mathscr{L}_{p}(K)(\Psi_{\mathbf{g}}(\ell)) &= \\ & \qquad \qquad \mathbf{f}(k, l) \cdot \mathscr{L}_{p}(\mathbf{f}/K)(k, \Psi_{\mathbf{g}h}(k, \ell)) \cdot \mathscr{L}_{p}(\mathbf{f}/K)(k, \Psi_{\mathbf{g}h'}(k, \ell)). \end{aligned}$$

*Proof.* We have a decomposition of classical L-functions given by

(17) 
$$L(f_k \otimes g_\ell \otimes h, \frac{k+\ell-1}{2}) = \mathcal{E}ul_N(k,\ell) \cdot L(f_k, \Psi_{gh}(k,\ell)^{-1}, 0) \cdot L(f_k, \Psi_{gh'}(k,\ell)^{-1}, 0).$$

Since  $p \nmid N$ , the Euler factor in  $\mathcal{E}ul_N(k,\ell)$  interpolates p-adically and gives rise to a two-variable  $L_0$ -admissible function. Combine equations (15) and (17) with the interpolation formula (3). Then use equations (6) and (9) to replace the classical L-functions with their respective p-adic avatar. An easy check shows that the periods simplify.

Define

$$\omega(k,\ell) := I(\breve{g}_{\ell}^*, \breve{f}_{k}^*, h^*)/I(\breve{g}_{\ell}, \breve{f}_{k}, \breve{h})$$

as in [DLR, (25)]. Recall also the local constants  $W_v$  for  $v \mid N\infty$  introduced in (4). Thanks to Lemma 5.3 one obtains the following equality, true for every  $(k, \ell) \in \mathcal{U}_{fq}^{\text{cl}}$ :

(18) 
$$\mathcal{L}_{p}^{g}(\check{\mathbf{f}}, \check{\mathbf{g}}, \check{h})(k, l)^{2} \cdot \mathcal{L}_{p}(K)(\Phi_{g}(\ell))^{2} = \mathbf{f}_{0}(k, \ell) \cdot \mathcal{L}_{p}(\mathbf{f}/K)(k, \Psi_{gh}(k, \ell))^{2} \cdot \mathcal{L}_{p}(\mathbf{f}/K)(k, \Psi_{gh'}(k, \ell))^{2}$$

where

$$\mathfrak{f}_0(k,\ell) = \frac{\mathcal{E}ul_N(k,\ell)}{\mathfrak{f}_{\mathrm{Pet}}(\ell)^2} \cdot \frac{\mathfrak{f}_{\mathrm{G}H}(\ell,k,1) \cdot \mathfrak{f}_{\mathrm{K}}(\Phi_g(\ell))^2}{\mathfrak{f}_{\mathrm{B}DP}(\Psi_{gh}(k,\ell)) \cdot \mathfrak{f}_{\mathrm{B}DP}(\Psi_{gh'}(k,\ell))} \cdot \frac{\prod_{v|N\infty} W_v}{\omega(k,\ell)}.$$

Let us show that  $\mathfrak{f}_0$  is  $L_0$ -admissible. First of all, it is immediate to verify that the powers or  $\pi$  appearing in the various fudge factors cancel out. Notice also that almost all terms appearing in the numerator and denominator of the definition of  $\mathfrak{f}_0(k,\ell)$  are  $L_0$ -admissible for obvious reasons, following directly from the definitions. This is the case except for the constants  $W_v$ ,  $\omega(f_k, \Psi_{qh})\omega(f_k, \Psi_{qh'})$  and  $\omega(k,\ell)$ , which need to be considered in more detail.

As for the local constants  $W_v$  is concerned, we already argued in §2 that for  $v \mid N$  the test vectors can be chosen in such a way that  $W_v$  lies in  $\mathbb{Q}^{\times}$  and do not depend on  $(k,\ell)$ . Moreover, the recent results of [Wo2] guarantee that in this setting  $W_{\infty} = 1$  for all  $(k,\ell)$ .

As for the global constant  $\omega(f_k, \Psi_{qh})\omega(f_k, \Psi_{qh'})$ , using lemma 5.4 one derives that:

$$\omega(f_k, \Psi_{gh})\omega(f_k, \Psi_{gh'}) = \frac{\psi_{g,\ell-1}(\mathfrak{N})^2 \psi_h(\mathfrak{N}) \psi_h'(\mathfrak{N})}{N_f^{\ell-1}}.$$

Since  $p \nmid N_f$ , it follows that the function  $\omega(f, \Psi_{gh})\omega(f, \Psi_{gh'})$  is  $L_0$ -admissible (in fact it is  $K(\psi_q, \psi_h)$ -admissible).

The  $L_0$ -admissibility of the function  $\omega(k, l)$  follows by the same argument as in the last part of the proof of [DLR, Theorem 3.9]. Hence we have proved that  $\mathfrak{f}_0$  is  $L_0$ -admissible and the theorem follows after taking the square-roots on both sides of (18).

Corollary 5.6 (Elliptic Stark conjecture). There exists a scalar  $\lambda \in L^{\times}$  such that:

$$\int_{\breve{\gamma}_{q\alpha}} \breve{f} \cdot \breve{g} = \lambda \cdot \frac{\log_{E,p}(P_{\psi_1'}) \log_{E,p}(P_{\psi_2'})}{\log_p(u_{\psi_0'})}.$$

*Proof.* Notice that we have

$$\Phi_g(1) = \psi_q'/\psi_g = \psi_0', \quad \Psi_{gh}(2,1) = \psi_1' \quad \text{and} \quad \Psi_{gh'}(2,1) = \psi_2'.$$

Using Theorem 5.5 and equations (5), (8) and (10) one obtains the desired result.

### 6. An explicit formula

Let us assume now that  $\psi = \psi_g$  is a ring class character of conductor  $c \in \mathbb{Z}_{\geq 1}$  and  $\psi_h = 1$ is the trivial character. This implies that  $h = \theta_{\psi_h} = E_{1,\chi_K}$  is the Eisenstein series associated with the pair of Dirichlet characters  $(1,\chi_K)$ . In this setting the characters defined in equations (13) and (14) coincide, i.e.  $\Psi_{gh} = \Psi_{gh'}$ , and thus we simply denote this character  $\Psi_g$ .

In this section we exploit the results obtained previously to prove an explicit formula for the scalar  $\lambda$  that makes its appearance in Theorem 1.1, in the particular scenario considered here. In order to do so we invoke yet another p-adic L-function in addition to the ones introduced so far. Namely, the Hida-Rankin p-adic L-function associated by Hida in Chapter 10 of [Hi93] to the convolution of two Hida families  $\mathbf{f}$ ,  $\mathbf{g}$  of cusp forms.

We follow the conventions and notations employed in [BDR1, §2] and [CR, §3], where Hida-Rankin's p-adic L-function  $\mathscr{L}_p{}^g(\mathbf{f},\mathbf{g})$  is described as a function on  $\mathcal{U}_{fg}$  and is characterized by the following interpolation property: for every  $(k,\ell) \in \mathcal{U}_{fg}^{\text{cl}}$  such that  $\ell \geq k+1$ , it holds that

(19) 
$$\mathscr{L}_{p}^{g}(\mathbf{f}, \mathbf{g})(k, \ell) = \mathfrak{e}_{HR}(k, \ell) \cdot \mathfrak{a}_{HR}(k, \ell) \cdot \mathfrak{f}_{HR}(k, \ell) \cdot \frac{L(\mathbf{f}_{k} \otimes \mathbf{g}_{\ell}, (\ell + k - 1)/2)}{\langle \mathbf{g}_{\ell}, \mathbf{g}_{\ell} \rangle_{N}}$$

where

- $\mathfrak{e}_{HR}(k,\ell) = \mathfrak{e}_{GH}(k,\ell,1);$
- $\mathfrak{a}_{HR}(k,\ell) = N/\sqrt{D_K}$   $\mathfrak{f}_{HR}(k,\ell) = (-1)^{\frac{\ell-k-1}{2}} \cdot \left(\frac{\ell-k-1}{2}\right)! \left(\frac{\ell-k-3}{2}\right)! \cdot 2^{-2\ell+1} \pi^{-\ell}.$

Moreover, by [CR, Proposition 3.2] we dispose of the following formula for the value of  $\mathcal{L}_p^g(\mathbf{f},\mathbf{g})$  at the point  $(k,\ell)=(2,1)$  lying outside the region of interpolation. Keeping the notations of the introduction and section 2, we have:

(20) 
$$\mathscr{L}_{p}^{g}(\mathbf{f},\mathbf{g})(2,1) = \int_{\check{\gamma}_{g_{\alpha}}} \check{f} \cdot \check{h}.$$

Recall the number field appearing in the statement of Theorem 5.5.

**Proposition 6.1.** There exists a two-variable L-admissible function  $f(k,\ell)$  such that the following factorization of p-adic L-functions holds:

$$\mathscr{L}_p^g(\mathbf{f}, \mathbf{g})(k, \ell) \cdot \mathscr{L}_p(K)(\Phi_g(\ell)) = \mathfrak{f}(k, \ell) \cdot \mathscr{L}_p(\mathbf{f}/K)(\Psi_{gh}(k, \ell))^2.$$

*Proof.* After setting  $\Psi := \Psi_{ah}(k,\ell)$ , it is easy to verify that

$$L(\check{g}_{\ell} \otimes \check{f}_{k}, (\ell+k-1)/2) = \mathcal{E}ul_{N}(k,\ell) \cdot L(f_{k}, \Psi, 0)$$

where  $\mathcal{E}ul_N(k,\ell)$  stands for a product of Euler factors at primes dividing N. The function  $\mathcal{E}ul_N(k,\ell)$  gives rise to an admidsible function as shown in [CR, Lemma 4.1]. From this, the proof proceeds along very similar lines as in the proof of Theorem 5.5 and [CR, Theorem 4.4], replacing the interpolation formula of the Garrett-Hida triple-product L-function by the one provided above.

We content here to record an explicit expression for the admissible function appearing in the above statement, namely

(21) 
$$\mathfrak{f}(k,\ell) = \mathcal{E}ul_N(k,\ell) \cdot (-1)^{\frac{k}{2}} \frac{\Im(D_K c^2)}{\Im(N)} \frac{N \cdot 2^{-\sharp\{q|(D_K,N_f)\}}}{h_c \cdot D_K} \frac{\psi_{\ell-1}(\mathfrak{N})}{c^{\ell+1} N_f^{\frac{\ell-1}{2}}}$$

where for any integer  $M \geq 1$  we set  $\Im(M) := [\operatorname{SL}_2(\mathbb{Z}) : \Gamma_0(M)] = \prod_{q^{n_q} \mid \mid M} q^{n_q - 1} (q + 1).$ 

Corollary 6.2. Assume  $N_f = D_K$  and  $g = h = \text{Eis}_1(1, \chi_K)$ . Then

$$\int_{\gamma_{q_{\alpha}}} f \cdot h = \lambda \cdot \frac{\log_{E,p}^{2}(P_{\psi'})}{\log_{p}(u_{\psi'}^{2})}$$

with

$$\lambda = \frac{(p - a_p(f)\psi(\wp))^2}{p \cdot c_E^2} \cdot \frac{\lambda_0}{h_K g_K}$$

where:

$$\lambda_0 = \begin{cases} \frac{1}{p-1} & \text{if } \psi^2 = 1, \text{ that is to say, if } g \text{ is Eisenstein} \\ \frac{12}{p-(p+1)\psi^2(\bar{\wp})+\psi^4(\bar{\wp})} & \text{if } \psi^2 \neq 1, \text{ that is to say, if } g \text{ is cuspidal.} \end{cases}$$

*Proof.* The value of the admissible function f appearing in (21) at the point  $(k, \ell) = (2, 1)$  is

$$\mathfrak{f}(2,1) = -\frac{2\psi(\mathfrak{N})}{h_K g_K}.$$

The result now follows upon combining this with formulae (8), (10) and (20).

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